Table I. Magnetic Susceptibilities<sup>a,b</sup>

Compound	Color	10 <sup>6</sup> χ <sub>m</sub> , cgs	$\mu_{eff}$ (obsd)	Theor <sup>13</sup>
Yb	Orange	7,800	4.31	4.5
Er	Pink	39,000	9.64	9.6
Gd	Yellow	26,880	7.98	7.94

<sup>a</sup> Measured by the Faraday method at 300°K. <sup>b</sup> Estimated error is  $\pm 3\%$ ; values are corrected for diamagnetism.

dienyl<sup>5</sup> and cyclooctatetraenyl<sup>6</sup> ligands. The only wellcharacterized  $\sigma$ -bonded derivatives presently known are the Yb and Lu tetrakis(2.6-dimethylphenyl) derivatives, isolated as lithium salts.<sup>7</sup> We wish to report the preparation of a new class of organolanthanide compounds,  $(\eta^5 - C_5 H_5)_2 Ln C \equiv CPh (Ph = phenyl; Ln = Gd,$ Er, and Yb). The Gd and Er derivatives represent the first examples of compounds containing a  $\sigma$ -bonded organic moiety bound to the respective metals.

These complexes are formed in good yields by the reaction of the particular  $(\eta^5-C_5H_5)_2LnCl^8$  and lithium phenylacetylide<sup>9</sup> in THF at  $-20^{\circ}$ . The compounds are purified by continuous extraction into benzene.<sup>10</sup> The compounds are destroyed by oxygen (though the Gd compound is stable in air for a few minutes) and both the phenylacetylide and cyclopentadienyl ligands are hydrolyzed by  $H_2O$ .

In the infrared spectra, all three compounds exhibit strong sharp C=C stretches at  $\sim 2050 \text{ cm}^{-1.11}$  In addition, strong bands at 692 and 760 cm<sup>-1</sup> are characteristic of a monosubstituted phenyl group. The bands at 3100, 1450, 1010, and 775  $cm^{-1}$  are indicative of a  $\eta^5$ -bonded cyclopentadienyl moiety.<sup>12</sup>

The magnetic susceptibilities of the compounds are reported in Table I, and are indicative of the metal in the tripositive state.13 One would not expect to see much deviation from these values as the 4f orbitals seem to be quite well shielded and should be perturbed only slightly by ligand fields.<sup>14</sup>

Visible spectra were taken for all of the complexes. The spectra of the erbium compound is interesting as it differs somewhat from the "free" Er<sup>3+</sup> ion. The bands are broadened slightly and the values for  $\epsilon$  are larger by over a factor of ten.<sup>15,16</sup> Bukietynska and Choppin<sup>18</sup>

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in THF.

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(11) In contrast, the C  $\equiv$  C stretch for ethynylbenzene comes at  $\sim$  2085 cm<sup>-1</sup>.

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suggested that changes of the intensity and shape of the hypersensitive<sup>19</sup> absorption band are due to the formation of inner-sphere complexes to a limited extent. High values can be related to the predominant existence of species in solution with low symmetries<sup>20</sup> and/or some degree of covalency. Since the symmetries in the starting material and products should be essentially identical, the observed enhancement lends support to these complexes being more covalent in nature. Also, a large charge transfer band is present which originates in the uv and extends into the visible region. This indicates that there is at least some interaction. other than strictly ionic, between the ligand and the

metal. The origin of the charge transfer is probably ligand to metal.<sup>6,21</sup> However, this interaction need not be a large one to be observed. Also, it is not clear whether the electron is transferred into a 4f or 5d orbital

The visible spectra of the Yb complex is uninformative as it is dominated by a large charge transfer band which is present in the  $(\eta^5 - C_5 H_5)_2$ YbCl compound.<sup>22</sup> The Gd complex (an f<sup>7</sup> compound) is devoid of all transitions except for the charge transfer band as seen in the Er complex.

Though heretofore unknown for most lanthanides. the formation of metal-carbon  $\sigma$  bonds by our reaction scheme appears to be a general one and should work for other ligands and metals as well.<sup>23,24</sup> Further studies are in progress to elucidate the nature of the lanthanide–carbon  $\sigma$  bond.

Acknowledgment. We wish to thank the United States Air Force Office of Scientific Research, AFOSR 71-2017, for financial support and Dr. Allen Gebala for helpful discussions.

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(22) This is the only lanthanide whose cyclopentadienide complexes do not correspond closely to the free ion.

(23) Preliminary results in our laboratories indicates this to be so,

(24) F. A. Hart, et al., postulate<sup>7</sup> that lanthanide-carbon  $\sigma$  bonds are most favored for the heavier elements of this series. Maginn, et al., also note<sup>8</sup> that they were only able to prepare  $(\eta^5-C_5H_5)_2LnCl$  for the heavier lanthanides and that this was due to some subtle effect of the lanthanide contraction.

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## Transverse Relaxation in Homonuclear **Coupled Spin Systems**

Sir

It is often stated that for nonviscous liquids, nuclear spin-spin relaxation times  $(T_2)$  are equal to spin-lattice relaxation times  $(T_1)$  unless "slow" processes occur. This notion, which is based on an unwarranted extension of the Bloch equations, is incorrect for the general case of molecules with scalar coupling. Data presented here, together with a complete density matrix analysis,<sup>1</sup> show that the presence of homonuclear scalar coupling gives rise to  $T_2$  values which are consid-

(1) R. L. Vold and R. R. Vold, to be submitted for publication.

erably shorter than corresponding  $T_1$ 's, even in the absence of slow processes such as chemical exchange and relaxation by scalar coupling to rapidly relaxing quadrupolar nuclei. Similar effects of heteronuclear coupling to other spin 1/2 nuclei have been described previously. 2-4

Table I shows relaxation times measured for selected

Table I. Proton Relaxation Times and (Rates)<sup>a</sup>

	$T_1(R_I)$	$T_2(R_2)$	${ar T}_2~({ar R}_2)^b$	
	Isopror	yl Iodide		
$(CH_3)_2^c$	7.3(0.137)	7.0 (0.143)	7.8 (0.128)	
CHd	13.7 (0.073)	7.9 (0.126)		
	Isoprop	yl Acetate		
$(CH_3)_{2^{c}}$	6.2(0.161)	6.2 (0.161)		
CHd	13.2 (0.076)	7.1 (0.141)	6,7 (0,149)	
CH <sub>3</sub> CO	8.5 (0.118)	7.8 (0.128)		

<sup>a</sup>  $R_1 = 1/T_1$  and  $R_2 = 1/T_2$  sec<sup>-1</sup>. <sup>b</sup> Weighted average relaxation times calculated from eq 1. c Equal values were obtained for both lines of the doublet within experimental error. <sup>d</sup> All components of the septet relaxed at the same rate. Listed values were obtained from the center line. Limited resolution (nonspinning sample,  $\Delta \nu_{1/2} \sim 2$  Hz) prevented detection of small differences. • A small deviation from single exponential behavior was observed; the reported  $T_1$  was obtained from the initial slope. For a justification of this procedure, see ref 16. If At  $2\tau = 1.2$  msec,  $T_2$  increased to 8.4 sec. Spin echo spectra (R. L. Vold and R. R. Vold, J. Magn. Resonance, 13, 38 (1974)) show that long range coupling to the acetyl protons is less than 0.05 Hz. The short value of  $T_2$  at  $2\tau =$ 200  $\mu$ sec may be due to a nonzero value for this coupling.

proton lines of neat isopropyl iodide and isopropyl acetate. All measurements were performed at 26° on a 55-MHz spectrometer described elsewhere.<sup>5,6</sup> The samples were degassed by distillation at  $\sim 10^{-5}$  Torr followed by several freeze-pump-thaw cycles.  $T_1$ values were obtained by the usual  $180^{\circ}-\tau-90^{\circ}-FT$  sequence and  $T_2$  values by a similar Fourier transform version7 of the Carr-Purcell-Meiboom-Gill (CPMG) spin echo sequence with a spacing  $2\tau = 200 \ \mu sec$  between refocusing pulses. Semilogarithmic plots spanning  $\sim$ 1.5 decades were linear except where noted, and the values of  $T_1$  and  $T_2$  are considered accurate<sup>8</sup> within ~3%.

At first glance the short  $T_2$  observed for the methine proton in isopropyl iodide might be ascribed to scalar relaxation by coupling to the rapidly relaxing iodine. However, the difference,  $1/T_2 - 1/T_1 = 0.053 \text{ sec}^{-1}$  for this proton is ca. eight times larger than that found for methyl iodide,<sup>4</sup> and a similar large difference, 0.075  $sec^{-1}$ , is observed for the methine proton in isopropyl acetate. Relaxation in isopropyl acetate is certainly dominated by dipole-dipole interactions, and infrared,

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microwave, and ultrasonic measurements on esters<sup>9</sup> have demonstrated that internal rotation rates are much too high to affect  $T_2$ 's by chemical exchange.

The differences between  $T_1$  and  $T_2$  values are due to the combined effects of scalar coupling and the rf pulses on the motion of the transverse magnetization vectors during the spin echo decay. Very small pulse intervals were used to avoid modulation of the spin echo decay with consequent phase and amplitude distortion of the spectra, <sup>10,11</sup> as well as other errors due to field instability and inhomogeneity.<sup>5,12</sup> In this "spin lock" limit, the magnetizations are refocused before any appreciable differential precession occurs, and the observed transverse decay times may also be designated  $T_{1_0}$ <sup>13,14</sup> In a CPMG spin echo experiment spin locking occurs if the pulse rate  $1/2\tau$  is much larger than the frequency spread of the spectrum.

The following heuristic argument explains the short  $T_2$  values observed under spin lock conditions. Chemical shifts are all effectively zero for the duration of the CPMG pulse train, while the scalar coupling is unaltered.<sup>15</sup> Therefore the effective coupling strength is extremely large, and transverse relaxation of any line should be indistinguishable from that of any other line. Thus, if a set of spins,  $A_n X_m$  with  $J_{AX} = 0$ , has transverse relaxation times  $T_{2A^0} = T_{1A}$  and  $T_{2A^0} = T_{1X}$  in the spin lock limit, then all lines in the scalar coupled system  $(J_{AX} \neq 0)$  will relax with the same transverse relaxation time  $\overline{T}_2$  where

$$1/\overline{T}_2 = (n/T_{2A^0} + m/T_{2X^0})/(n+m)$$
(1)

If we apply eq 1 to the  $T_1$  values listed in Table I, we obtain calculated values of  $\overline{T}_2$  listed in the third column. Agreement with observed  $T_2$  values is satisfactory considering the simplicity of the argument and demonstrates that scalar relaxation by coupling to iodine is unimportant for isopropyl iodide.

Complete density matrix calculations on AX and  $AX_2$ spin systems<sup>1</sup> confirm the essential features of the argument leading to eq 1. Equation 1 neglects the fact that scalar coupling does not mix states of different symmetry and that different relaxation mechanisms have different symmetry properties. Differences between observed and calculated  $T_2$ 's are due to effects of this nature. For sufficiently simple spin systems it is possible to experimentally distinguish contributions of different relaxation mechanisms on this basis.

It is important to note that complexities inherent in the density matrix analysis cannot be avoided by using long pulse intervals, thereby allowing free precession to develop. Under ideal conditions  $T_2$ 's observed as  $2\tau \rightarrow \infty$  will equal those obtained from line widths in a perfectly homogeneous field. They will normally be less than  $T_1$  due to irreversible dephasing from spin-lattice relaxation combined with scalar coupling.6.16

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## Efficiency of Excited State Carbonyl Production from the Thermolysis of 3,3-Diphenyl-1,2-dioxetane

Sir:

The unique ability of 1,2-dioxetanes to generate electronically excited carbonyl molecules during thermal decomposition has attracted considerable interest recently in these peroxides.<sup>1</sup> Thermochemical calculations suggest that sufficient energy is available from the thermal decomposition of 1,2-dioxetanes to obtain one of the carbonyl products in an excited state<sup>1e</sup> (efficiency  $= \alpha = 1.0$ ). However, experimental  $\alpha$  values range from 0.005 to about 1.0.11, j, n.o., s, z, aa, bb It is of considerable interest to determine what factors are responsible for the efficient conversion of the available thermal energy into electronic energy.

With this longer range goal as our objective, we now report the efficiency of excited state carbonyl production from the thermolysis of 3,3-diphenyl-1,2-dioxetane (DPD). Isomerization of trans-stilbene was used to determine the total efficiency of excited state carbonyl formation<sup>1z</sup> from DPD, while excited state benzophenone molecules were specifically monitored by cycloaddition with 2-methyl-2-butene.

In Scheme I, the pertinent processes are shown for the decomposition of DPD in the presence of *trans*-stilbene (t-S) where *cis*-stilbene (c-S) is produced. The total excited state carbonyl molecules (benzophenone and formaldehyde) are represented by K\* and the ground state species by K. The apparent quantum yield  $(\Phi_{App})$  for isomerization of t-S is given by eq 1, where

$$\Phi_{App} = \alpha \Phi_{ET} \Phi_{t \to c} \tag{1}$$

 $\Phi_{t\rightarrow c}$  is the quantum yield for isometization of t-S and

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Scheme I

$$O - O$$

$$(C_{6}H_{6})_{2}C - CH_{2} \xrightarrow{k_{D}} \alpha K^{*} + (1 - \alpha)K$$

$$DPD$$

$$K^{*} \xrightarrow{k_{d}} K$$

$$K^{*} + t \cdot S \xrightarrow{k_{q}} K + t \cdot S^{*}$$

$$t \cdot S^{*} \xrightarrow{k_{i}} c \cdot S$$

$$t \cdot S^{*} \xrightarrow{k_{d}'} t \cdot S$$

 $\Phi_{ET}$  is the quantum yield for energy transfer between K\* and *t*-S (eq 2).

$$\Phi_{\rm ET} = \frac{k_{\rm q}[t-{\rm S}]}{k_{\rm q}[t-{\rm S}] + k_{\rm d}}$$
(2)

DPD was allowed to decompose in degassed benzene solution with t-S at 45° through approximately ten half-lives.<sup>2</sup> The amount of c-S was determined by glc (3% SE-30 on Varaport-30, 5 ft  $\times$   $\frac{1}{8}$  in., 115°, flow 24 ml of N<sub>2</sub>/min, FID;  $t_r(min)$ , c-S (9.8), t-S (19)) and  $\Phi_{App}$  was calculated as the ratio of millimoles of c-S produced/millimoles of DPD decomposed. With  $\Phi_{t\rightarrow c}$ = 0.55<sup>3</sup> and assuming  $\Phi_{\rm ET}$  is unity,  $\alpha$  is calculated from eq 1. These data are given in Table I, where the con-

Table I. Isomerization of trans-Stilbene (t-S) from the Thermolysis of 3,3-Diphenyl-1,2-dioxetane (DPD) in Degassed Benzene at 45°

10³[DPD]₀, <i>M</i>	10 <sup>2</sup> [ <i>t</i> -S], <i>M</i>	10²Ф <sub>Арр</sub> ª	$10^2 \alpha^b$
1.98	2.00	1.77	3.22
19.8	2.50	1.85	3.36
5.10	2.50	1.72	3.13
5.10	5.00	2.16	3.93
1,98	10.0	1.89	3,44
19.8	20.0	2.01	3.65
		$Av 1.90 \pm 0.12$	Av 3.46 $\pm$ 0.23

<sup>*a*</sup>  $\Phi_{App}$  = millimoles of *c*-S produced/millimoles of DPD decom $b \alpha$  = total efficiency of excited state carbonyl production posed. from DPD, where  $\alpha = \Phi_{App}/0.55$  and  $\Phi_{t\rightarrow c} = 0.55$ ,  $\Phi_{ET} = 1.00$ .

centrations of both t-S and DPD are varied. The calculation of  $\alpha$  from eq 1, with the approximation that  $\Phi_{\rm ET}$  is unity, requires that  $k_{\rm q}[t-S] \gg k_{\rm d}$  in eq 2. This approximation appears reasonable for the data in Table I, where  $k_q$  is estimated to be 6.9  $\times$  10<sup>9</sup>  $M^{-1}$  sec<sup>-1</sup> at  $45^{\circ 4}$  and at the lowest *t*-S concentration  $(2.00 \times 10^{-2} M)$ ,  $k_{q}[t-S] = 1.4 \times 10^{8} \text{ sec}^{-1}$ . This value is then sufficiently large, compared to  $k_d \cong 10^5 \text{ sec}^{-1}$  at 23°<sup>8</sup> for benzophenone in benzene, to satisfy the approximation. As

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